The STELLA 3D Fokker-Planck Code

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History:

- Originally, calculation of f(v,θ,s_along_B;t) with v_{II}.df/ds-term (i.e. NOT bounce-averaged) used for neoclassical resistivity at all collisionalities (Sauter, Harvey, Hinton, Contr PI Phys (1994)), subsequently Sauter-bootstrap-current formulas (Sauter, Angioni, Lin-Liu, PoP (1999)). Part of CQL3D
- Then with E_{II}, applied to tokamak divertors (Kupfer, Harvey, Phys PI (1996)), called FPET
- Adjusted for multi-species application to solar wind (~2000), with Smirnov, renamed STELLA
- Recently B field variation added.

Divertor Application: Why Kinetic Modeling of SOL?

• Heat flux carried by electrons at $v \sim 3V_{Te}$ [Remember: $V_{coll} \propto 1/v^3$]

$$\Rightarrow \frac{\lambda_{mfp} |_{_{3V_{Te}}}}{L_{along B}} > 1 , typically$$

- ⇒ Kinetic effects important (often) for heat transport in SOL
 (Ad Hoc flux-limiting factor is often used in the fluid treatment)
- Localized cooling (e.g. Marf's) increases "temperature" gradient
 ⇒ Kinetic effects.

FPET/Stella

Solves Fokker-Planck Eqn for distr. function $f(v \text{[speed]}, \theta \text{[pitch-angle]})$, for e, i at each *s* [distance] along **B**:

$$\frac{\partial f(v,\theta,s,t)}{\partial t} + v_{\parallel} \frac{\partial f}{\partial s} + \frac{q E_{\parallel}}{m} \frac{\partial f}{\partial v_{\parallel}} = C(f) + D_{RF}(v,\theta,s) + S(v,\theta,s)$$

C(f) = Non-Linear RMJ Coulomb Collision Operator

 $D_{RF}(v, \theta, s)$, $S(v, \theta, s)$ are 3D RF diffusion terms and particle sources, with B.C.'s and E_{\parallel} for given flux(s), chosen to maintain charge neutrality.

Solution of finite difference eqns is by Alternating Direction Implicit: 1) Implicit vel. step based on CQL3D algorithm, with explicit s-dependence 2) Implicit (tri-diagonal) solve in s, with explicit vel-dependence

 E_{\parallel} is determined by an iterative procedure so that e,i-fluxes are equal. (Kupfer, Harvey, Phys. of Plasmas 3, 3644 (1996))

Boundary Conditions in FPET/Stella



Sequence of Calculations for Combined Electron-Ion Problem

(General Divertor Case)

- 0. Initialize with Maxwellian electron and ion profiles, plus sources.
- 1. Time step electrons to ~steady-state, with $E_{||k}$ chosen for constant $n_e(s)$.
- 2. Single time step ions with E_k determined (in 1.) from electrons. Ion Δt is much greater than electron Δt . This will permit ion density profile to relax.
- 3. Several electron time steps, to get quasi-state E_{\parallel}
- Time-step to steady-state
- Basically, electrons determine E_{\parallel} and ions accelerate to C_{s} due to E_{\parallel}

Comparison of Fluid and Kinetic Solutions in a Tokamak Divertor



Ion Distn is ~Shifted Maxwl at vel c_s + similar tail toward the divertor plate.



Modified heat conductivity Strong modification of ionization rates Etc

Another Application ECCD by Power Deposited in Localized EC Spot On A Field Line

ECH injected in narrow cone, heats localized spot on a range of flux surfaces. What keeps electron current going along **B**, outside of the spot?

> The connection length around the tokamak until the B-field connects back to the spot (2/3 overlap) is about 20 toroidal rotations.

L_{connection} ~ λ_{mfp}

Model $f_e(v,\theta,z;t)$ roughly in STELLA with periodic BCs and RF parallel (or perp) diffusion localized over 0.2 of total length of a field line, L~100 m, under typical tokamak plasma conditions.





Observations from Model for Localized ECH on Tokamak Flux Surface

- Electrons are on avg accelerated in the RF region, and collisionally slow down over the rest of the connection length, creating bunching (space charge) and E_{\parallel} .
- Pressure buildup is also an important part of the dynamics.
- Thus, a force, $F_e = -eE_{\parallel} \frac{1}{n_e} \frac{\partial p_e}{\partial z}$, in addition to momentum input from the RF. This force maintains continuity of the current within the toroidal circuit:
 - F_e : Retards the e-current in the RF region, accelerates it outside.
 - Net work around the circuit by Fe is zero: ∮ F_eds=0
 So, no net effect of these additional forces, relative to usual uniform plasma calcs of current drive efficiency.
- CD efficiency agrees with uniform plasma theory
- (Harvey et al., Proc. of RF Power in Plasmas Mtg, Annapolis (1999))
- The current turns on rather slowly:



• This model illustrates STELLA capabilities for calculation of the ambipolar electric field over a range of collisionality regimes, from collisionless to collisional.

Possible Relation of This Work to FRC

CQL3D Bounce Average Approach:

- May be adapted for electron tail heating by ECH or HHFW
- Trapped particles in field reversed region, transiting outside
- lons, could get some rough estimates, but gyro-orbits are large

STELLA non-bounce-averaged approach:

- Seems can be quite accurate for parallel electron transport, and has no limitations regards multiple minima of |B|
- Could work well for calculation of ambipolar electric fields, including the exhaust region. Would run the code on a radial array of flux surfaces.
- Sources and sink readily added for ionization and recombination
- RF diffusion heating can be added, coupling to GENRAY
- Ion modeling would not be very accurate, due to large gyro-orbits